

# Electrodynamics

SECTION

Electrostatics

LECTURE

Maxwell's equations

# Maxwell's equations

$$\nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} \rho$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{j}$$

- Charges are sources of the electric field
- The magnetic field has no sources (no monopoles)
- Time-dependent magnetic fields generate electric fields
- Time-dependent electric fields and currents generate magnetic fields

Vacuum permeability  $\mu_0 = 4\pi \cdot 10^{-7} \frac{\text{N}}{\text{A}^2}$

Vacuum permittivity  $\epsilon_0 = 8.854... \cdot 10^{-12} \frac{\text{As}}{\text{Vm}}$

# Maxwell's equations

Differential formulation

$$\nabla \cdot \mathbf{E} = \frac{1}{\epsilon_0} \rho$$

$$\nabla \times \mathbf{E} = -\cancel{\frac{\partial \mathbf{B}}{\partial t}}$$

Integral formulation

$$\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV$$

$$\oint_{\partial A} \mathbf{E} \cdot d\mathbf{l} = -\int_A \cancel{\frac{\partial \mathbf{B}}{\partial t}} \cdot d\mathbf{S}$$

# Electrodynamics

SECTION

Electrostatics

LECTURE

Poisson equation

# Poisson equation

- Maxwell  $\nabla \times \mathbf{E}(\mathbf{r}) = 0$
- We can introduce the electrostatic potential  $\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r})$

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = - \int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$

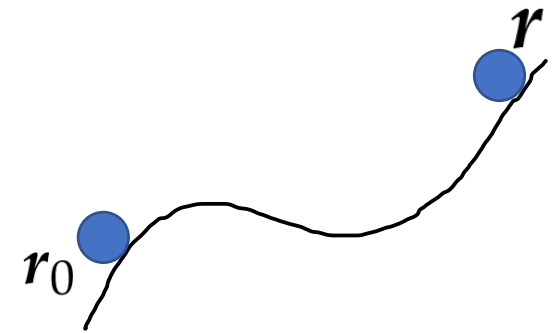
- Like in classical mechanics  $\mathbf{F}(\mathbf{r}) = -\nabla V(\mathbf{r})$

$$V(\mathbf{r}) - V(\mathbf{r}_0) = - \int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{F}(\mathbf{r}') \cdot d\mathbf{r}'$$

- Maxwell  $\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{1}{\epsilon_0} \rho(\mathbf{r})$

$$\Delta\varphi(\mathbf{r}) = -\frac{1}{\epsilon_0} \rho(\mathbf{r})$$

Poisson equation



# Electrodynamics

SECTION

Electrostatics

LECTURE

Rotational-symmetric  
charge distributions

# Poisson equation

- Maxwell  $\nabla \times \mathbf{E}(\mathbf{r}) = 0$
- We can introduce the electrostatic potential  $\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r})$

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = - \int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$

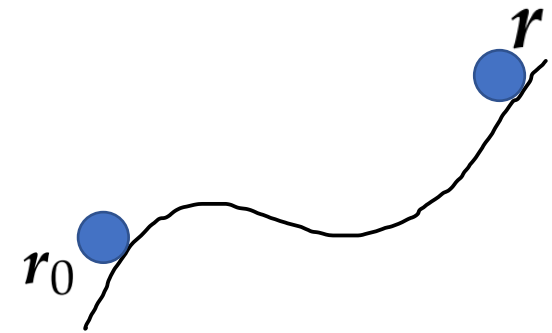
- Like in classical mechanics  $\mathbf{F}(\mathbf{r}) = -\nabla V(\mathbf{r})$

$$V(\mathbf{r}) - V(\mathbf{r}_0) = - \int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{F}(\mathbf{r}') \cdot d\mathbf{r}'$$

- Maxwell  $\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{1}{\epsilon_0} \rho(\mathbf{r})$

$$\Delta\varphi(\mathbf{r}) = -\frac{1}{\epsilon_0} \rho(\mathbf{r})$$

Poisson equation



# Coulomb's law

- Maxwell  $\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV = Q_v$

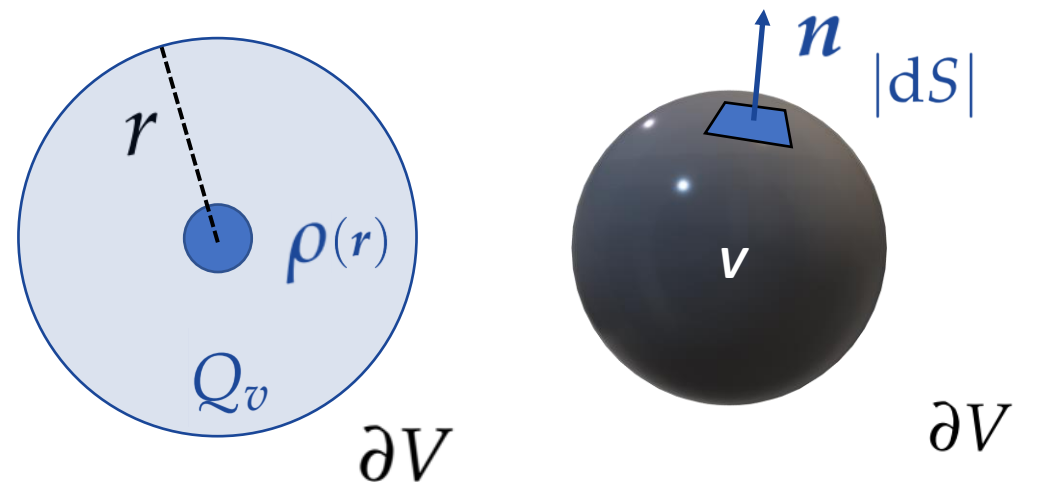
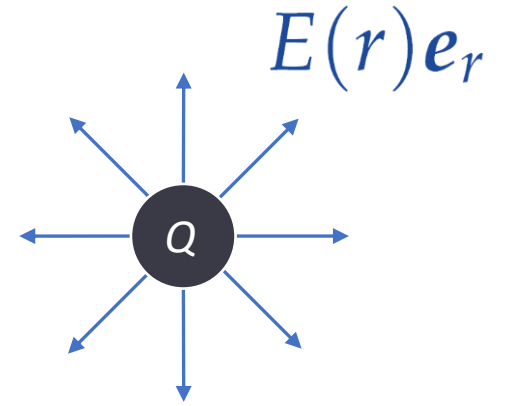
$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r \quad \mathbf{E} \parallel d\mathbf{S}$$

$$\epsilon_0 E(r) \oint_{\partial V} dS = Q_v$$

$$\epsilon_0 E(r) 4\pi r^2 = Q_v$$

$$E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2}$$

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2} \mathbf{e}_r$$



# Electrostatic potential

$$\Delta\varphi(\mathbf{r}) = -\frac{1}{\epsilon_0}\rho(\mathbf{r})$$

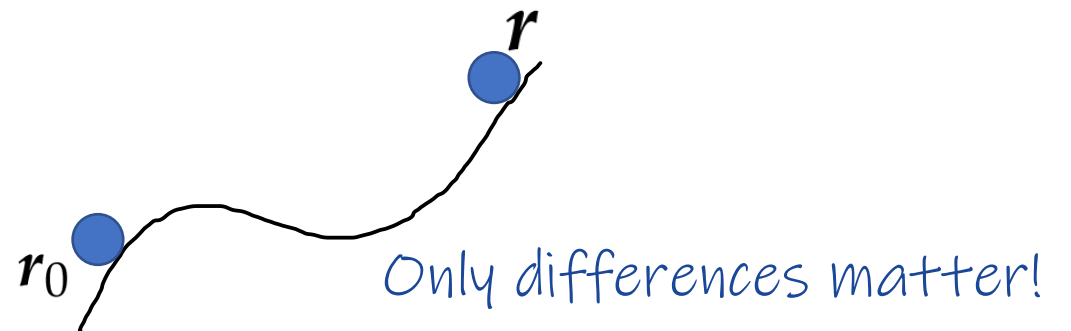
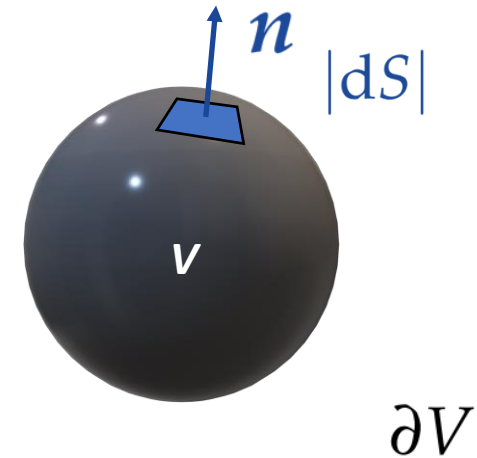
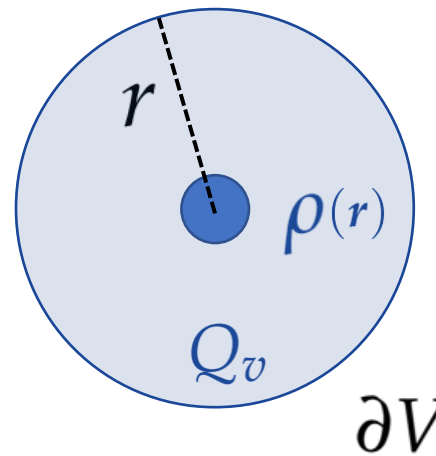
$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2} \mathbf{e}_r$$

If we have  $\mathbf{E}(\mathbf{r})$  it is easy

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = -\int_{r_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$

$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r} - \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r_0}$$

Solve Poisson equation?



# Electrodynamics

SECTION

Electrostatics

LECTURE

General charge  
distributions

# Electrostatic potential

$$\Delta\varphi(\mathbf{r}) = -\frac{1}{\epsilon_0}\rho(\mathbf{r})$$

Poisson equation

$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_1}{|\mathbf{r} - \mathbf{r}_1|}$$

- Point charge

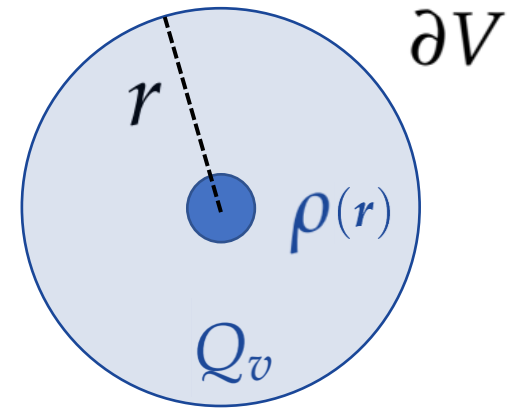
$$\varphi(\mathbf{r}) = \sum_i \frac{1}{4\pi\epsilon_0} \frac{Q_i}{|\mathbf{r} - \mathbf{r}_i|}$$

- Multiple point charges

- General charge distribution

$$\varphi(\mathbf{r}) = \int \frac{1}{4\pi\epsilon_0} \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'$$

$$Q = \int \rho(\mathbf{r}') d\mathbf{r}'$$



# Electrodynamics

SECTION

Electrostatics

LECTURE

Exercise:  
charged sphere

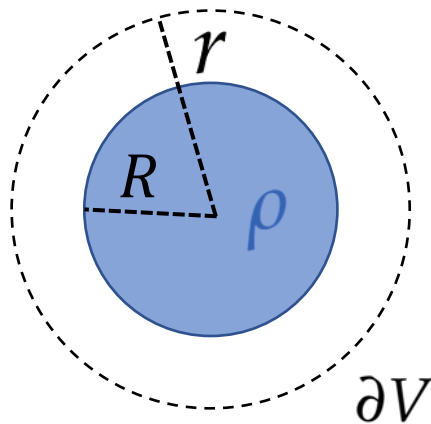
# Exercise

1. Calculate the electric field  $\mathbf{E}(\mathbf{r})$  of a homogeneously charged sphere (distinguish inside and outside).

*Hint: Use rotational symmetry to simplify the problem.*

*Hint: "Homogeneously charged" means that  $\rho = \text{const.}$*

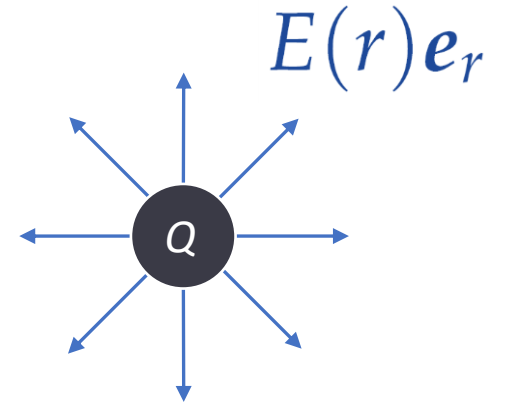
2. Calculate the electrostatic potential  $\varphi(\mathbf{r})$ .



# Coulomb's law

- Maxwell  $\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV = Q_v$

$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r \quad \mathbf{E} \parallel d\mathbf{S}$$



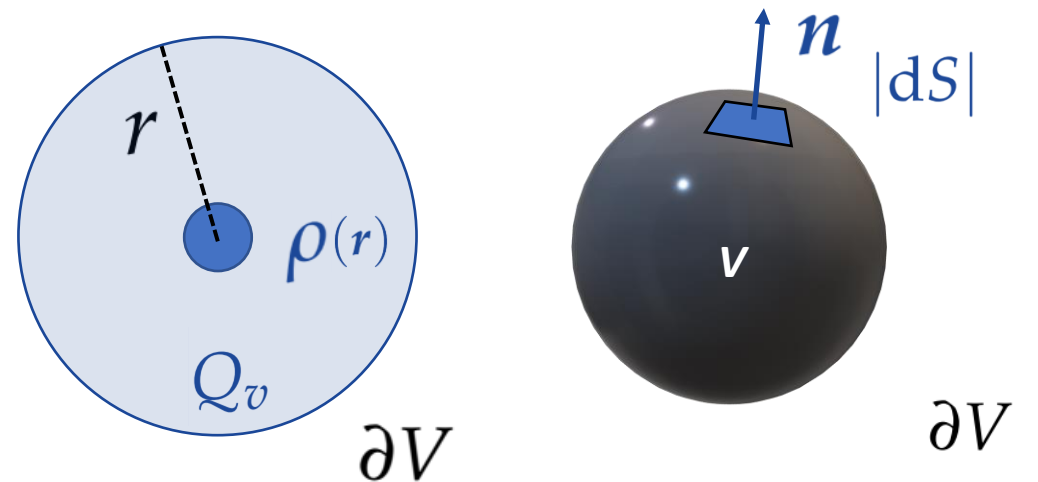
$$\epsilon_0 E(r) \oint_{\partial V} dS = Q_v$$

What about inside of the sphere?

$$\epsilon_0 E(r) 4\pi r^2 = Q_v$$

$$E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2}$$

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2} \mathbf{e}_r$$

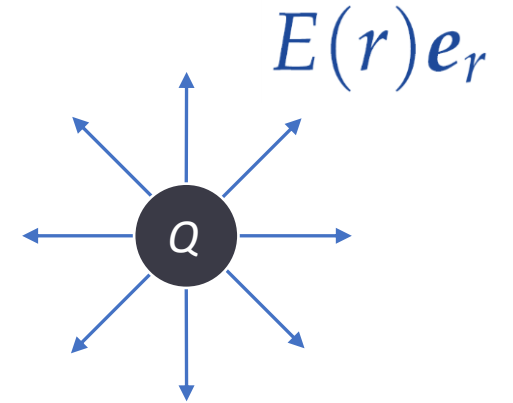


# Charged sphere

- Maxwell  $\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV$

$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r$$

$$\mathbf{E} \parallel d\mathbf{S}$$

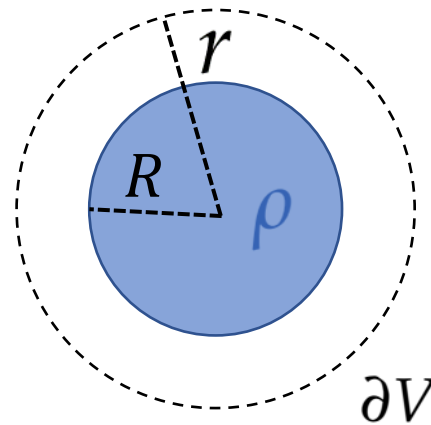


- Outside  $r > R$

$$\epsilon_0 E(r) \oint_{\partial V} dS = \frac{4}{3} \pi R^3 \rho = Q_v$$

$$\epsilon_0 E(r) 4\pi r^2 = \frac{4}{3} \pi R^3 \rho = Q_v$$

$$E(r) = \frac{R^3 \rho}{3\epsilon_0 r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2}$$



- Inside  $r < R$

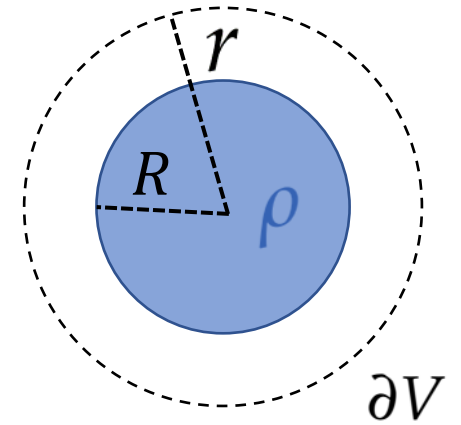
$$\epsilon_0 E(r) \oint_{\partial V} dS = \rho \int_V dV$$

$$\epsilon_0 E(r) 4\pi r^2 = \frac{4}{3} \pi r^3 \rho$$

$$E(r) = \frac{\rho}{3\epsilon_0} r$$

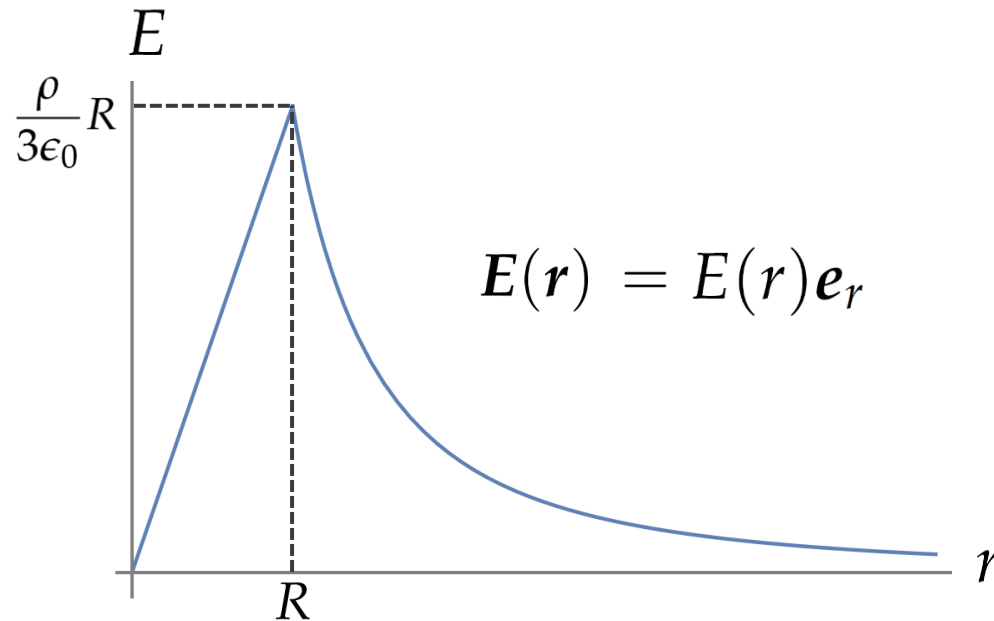
# Charged sphere

1. Calculate the electric field  $\mathbf{E}(r)$  of a homogeneously charged sphere (distinguish inside and outside).



- Inside  $r < R$

$$E(r) = \frac{\rho}{3\epsilon_0} r$$



- Outside  $r > R$

$$E(r) = \frac{R^3 \rho}{3\epsilon_0 r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2}$$

# Charged sphere

2. Calculate the electrostatic potential  $\varphi(\mathbf{r})$ .

- Inside  $r < R$

$$E(r) = \frac{\rho}{3\epsilon_0} r$$

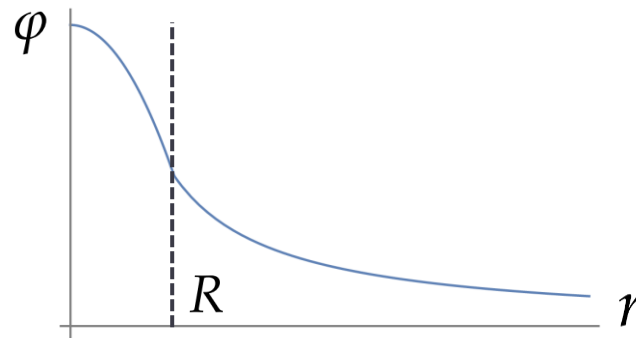
$$\varphi(r) = -\frac{\rho}{6\epsilon_0} r^2 + \frac{1\rho}{2\epsilon_0} R^2$$

$$\varphi(r) = \frac{\rho}{6\epsilon_0} (3R^2 - r^2)$$

- electrostatic potential

$$\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r})$$

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = -\int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$



- Outside  $r > R$

$$E(r) = \frac{R^3 \rho}{3\epsilon_0 r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r^2}$$

$$\varphi(r) = \frac{R^3 \rho}{3\epsilon_0 r} = \frac{1}{4\pi\epsilon_0} \frac{Q_v}{r}$$

# Electrodynamics

SECTION

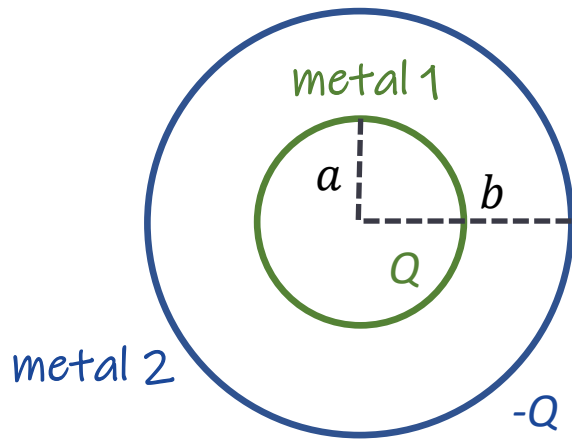
Electrostatics

LECTURE

Exercise:  
spherical capacitor

# Spherical capacitor

1. Calculate the electric field  $\mathbf{E}(\mathbf{r})$  of the following capacitor for all regions.
2. Calculate the electrostatic potential  $\varphi(\mathbf{r})$  for all regions.
3. Calculate the voltage between the two metals  $U = \varphi(r = a) - \varphi(r = b)$ .
4. Calculate the capacitance  $C = Q/U$ .



• Maxwell  $\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV$

$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r$$

$$\mathbf{E} \parallel d\mathbf{S}$$

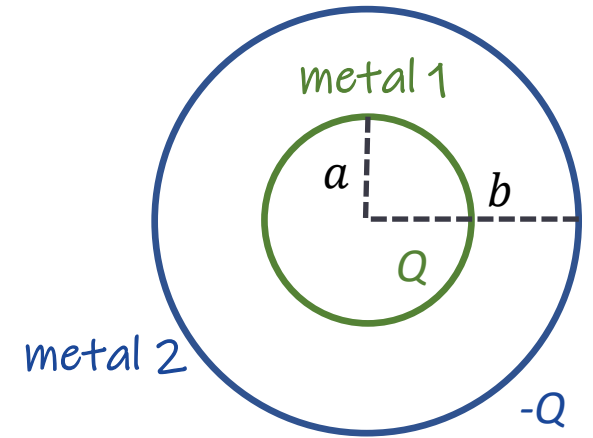
# Spherical capacitor

1. Calculate the electric field  $\mathbf{E}(\mathbf{r})$  of the following capacitor for all regions.

- Inside  $r < a$ :  $E(\mathbf{r}) = 0$

No enclosed charge  $\rho(\mathbf{r}) = 0$

$$\epsilon_0 \oint_{\partial V} \mathbf{E} \cdot d\mathbf{S} = \int_V \rho dV$$



- Between  $a < r < b$ :

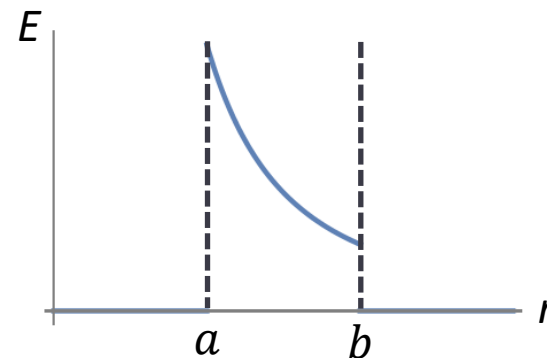
$$\epsilon_0 \oint_{\partial V} E dS = Q$$

$$\epsilon_0 4\pi r^2 E = Q$$

$$E(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \quad \mathbf{E}(\mathbf{r}) = E(\mathbf{r})\mathbf{e}_r$$

- Outside  $r > b$ :  $E(\mathbf{r}) = 0$

No enclosed charge  $Q_v = Q - Q = 0$



# Spherical capacitor

2. Calculate the electrostatic potential  $\varphi(\mathbf{r})$ .

- Inside  $r < a$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = \text{const.}$$

- Between  $a < r < b$ :  $E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$

$$\varphi(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r} + \text{const.}$$

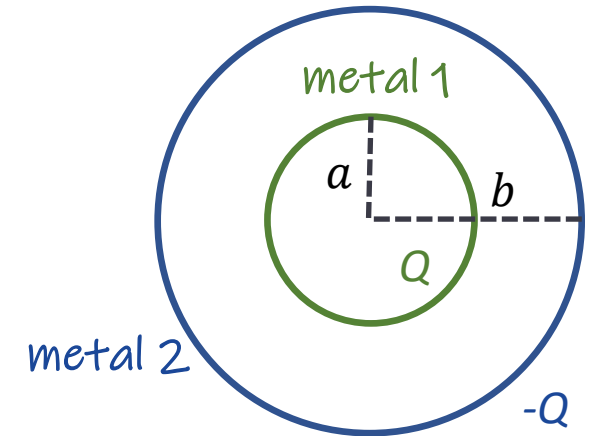
- Outside  $r > b$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = \text{const.}$$

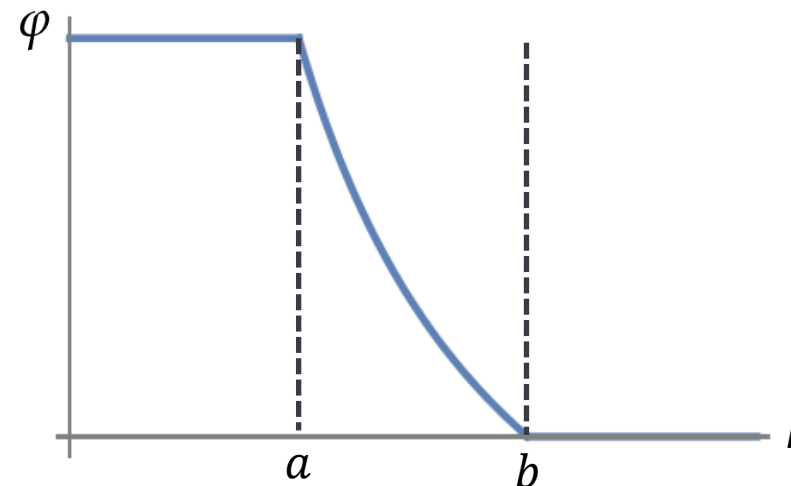
electrostatic potential

$$\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r})$$

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = -\int_{r_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$



$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r$$



# Spherical capacitor

2. Calculate the electrostatic potential  $\varphi(\mathbf{r})$ .

- Inside  $r < a$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right)$$

- Between  $a < r < b$ :  $E(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{1}{b} \right)$$

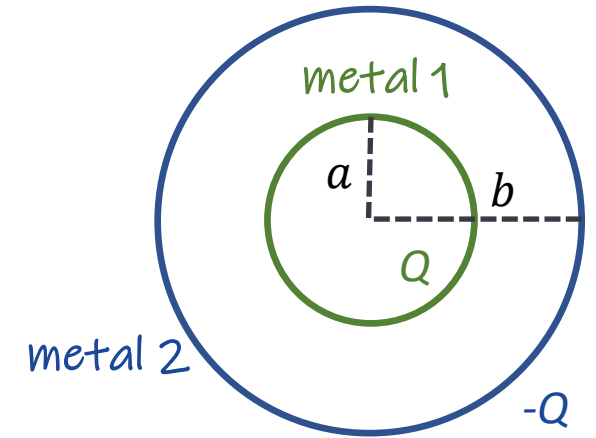
- Outside  $r > b$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = 0$$

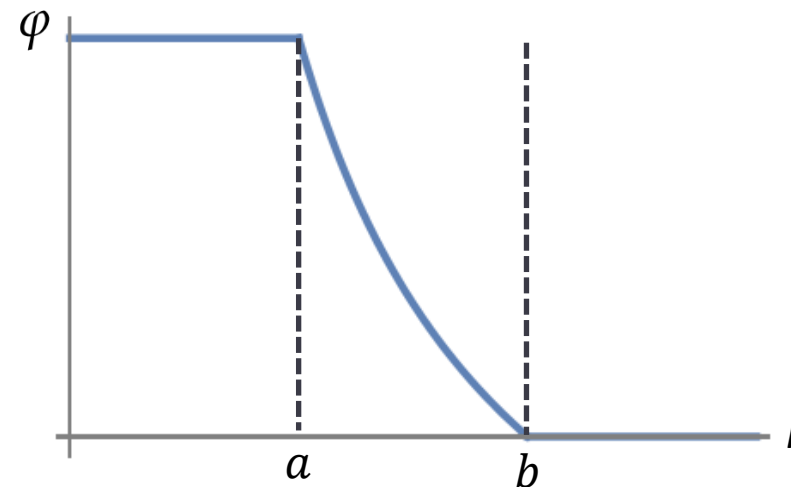
electrostatic potential

$$\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r})$$

$$\varphi(\mathbf{r}) - \varphi(\mathbf{r}_0) = -\int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{E}(\mathbf{r}') \cdot d\mathbf{r}'$$



$$\mathbf{E}(\mathbf{r}) = E(r)\mathbf{e}_r$$



# Spherical capacitor

3. Calculate the voltage between the two metals  $U = \varphi(r = a) - \varphi(r = b)$ .

• Inside  $r < a$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right)$$

$$U = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{b} \right)$$

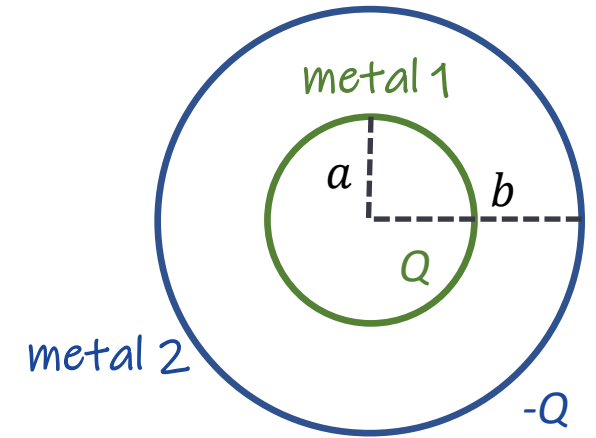
• Outside  $r > b$ :  $E(\mathbf{r}) = 0$

$$\varphi(\mathbf{r}) = 0$$

4. Calculate the capacitance  $C = Q/U$ .

$$C = Q/U = \frac{Q}{\frac{1}{4\pi\epsilon_0} Q \left( \frac{1}{a} - \frac{1}{b} \right)}$$

$$C = 4\pi\epsilon_0 \frac{ab}{b - a}$$



# Electrodynamics

SECTION

Electrostatics

LECTURE

Electric dipole

# Electric dipole

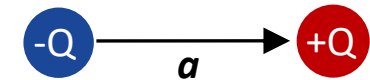
- For multiple point charges

$$\varphi(\mathbf{r}) = \sum_i \frac{1}{4\pi\epsilon_0} \frac{Q_i}{|\mathbf{r} - \mathbf{r}_i|}$$

- For electric dipole

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{1}{|\mathbf{r} + \mathbf{a}|} \right)$$

$$\mathbf{p} = Q\mathbf{a}$$



Taylor expansion

$$\frac{1}{|\mathbf{r} + \mathbf{a}|} = \sum_{n=0}^{\infty} \frac{1}{n!} (\mathbf{a} \cdot \nabla)^n \frac{1}{r}$$

For  $r \gg a$

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{1}{r} - \mathbf{a} \cdot \nabla \frac{1}{r} \right)$$

$$\nabla \frac{1}{r} = -\frac{\mathbf{r}}{r^3}$$

$$\nabla \frac{1}{r} = \nabla (x^2 + y^2 + z^2)^{-1/2} = \left( -\frac{1}{2} \right) (x^2 + y^2 + z^2)^{-3/2} \begin{pmatrix} 2x \\ 2y \\ 2z \end{pmatrix}$$

- Electrostatic potential

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \frac{\mathbf{a} \cdot \mathbf{r}}{r^3}$$

$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \mathbf{r}}{r^3}$$

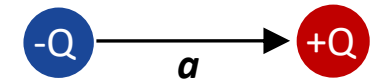
# Electric dipole

- Electrostatic potential  $\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{1}{|\mathbf{r} + \mathbf{a}|} \right)$

For  $r \gg a$

$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \mathbf{r}}{r^3}$$

$$\mathbf{p} = Q\mathbf{a}$$



- Electric field  $\mathbf{E}(\mathbf{r}) = -\nabla\varphi(\mathbf{r}) = -\frac{1}{4\pi\epsilon_0} \nabla \frac{\mathbf{p} \cdot \mathbf{r}}{r^3}$

$$\nabla \frac{\mathbf{p} \cdot \mathbf{r}}{r^3} = \nabla \frac{p_x x + p_y y + p_z z}{r^3} = \begin{pmatrix} p_x/r^3 \\ p_y/r^3 \\ p_z/r^3 \end{pmatrix} + (p_x x + p_y y + p_z z) \nabla \frac{1}{r^3}$$

$$\nabla \frac{\mathbf{p} \cdot \mathbf{r}}{r^3} = \frac{\mathbf{p}}{r^3} + \mathbf{p} \cdot \mathbf{r} \nabla \frac{1}{r^3}$$

$$\nabla \frac{1}{r^3} = \nabla (x^2 + y^2 + z^2)^{-3/2} = \begin{pmatrix} -3 \\ -2 \end{pmatrix} (x^2 + y^2 + z^2)^{-5/2} \begin{pmatrix} 2x \\ 2y \\ 2z \end{pmatrix}$$

$$\nabla \frac{1}{r^3} = -3 \frac{\mathbf{r}}{r^5}$$

$$\mathbf{E}(\mathbf{r}) = -\frac{1}{4\pi\epsilon_0} \nabla \frac{\mathbf{p} \cdot \mathbf{r}}{r^3} = -\frac{1}{4\pi\epsilon_0} \left( \frac{\mathbf{p}}{r^3} - 3\mathbf{p} \cdot \mathbf{r} \frac{\mathbf{r}}{r^5} \right) = -\frac{1}{4\pi\epsilon_0} \frac{r^2 \mathbf{p} - 3\mathbf{r}(\mathbf{p} \cdot \mathbf{r})}{r^5}$$

# Electric dipole

- Electrostatic potential

$$\varphi(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{1}{|\mathbf{r} + \mathbf{a}|} \right)$$

For  $r \gg a$

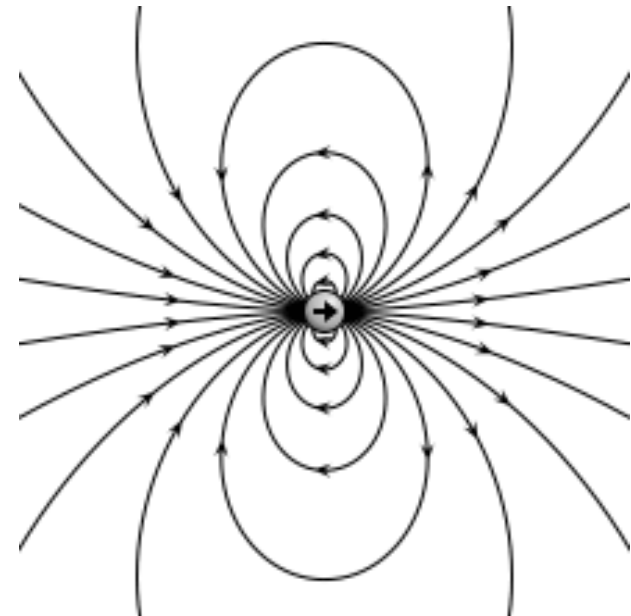
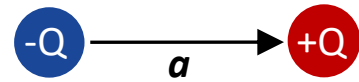
$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \mathbf{r}}{r^3}$$

- Electric field

$$\mathbf{E}(\mathbf{r}) = -\frac{1}{4\pi\epsilon_0} \frac{r^2 \mathbf{p} - 3\mathbf{r}(\mathbf{p} \cdot \mathbf{r})}{r^5}$$

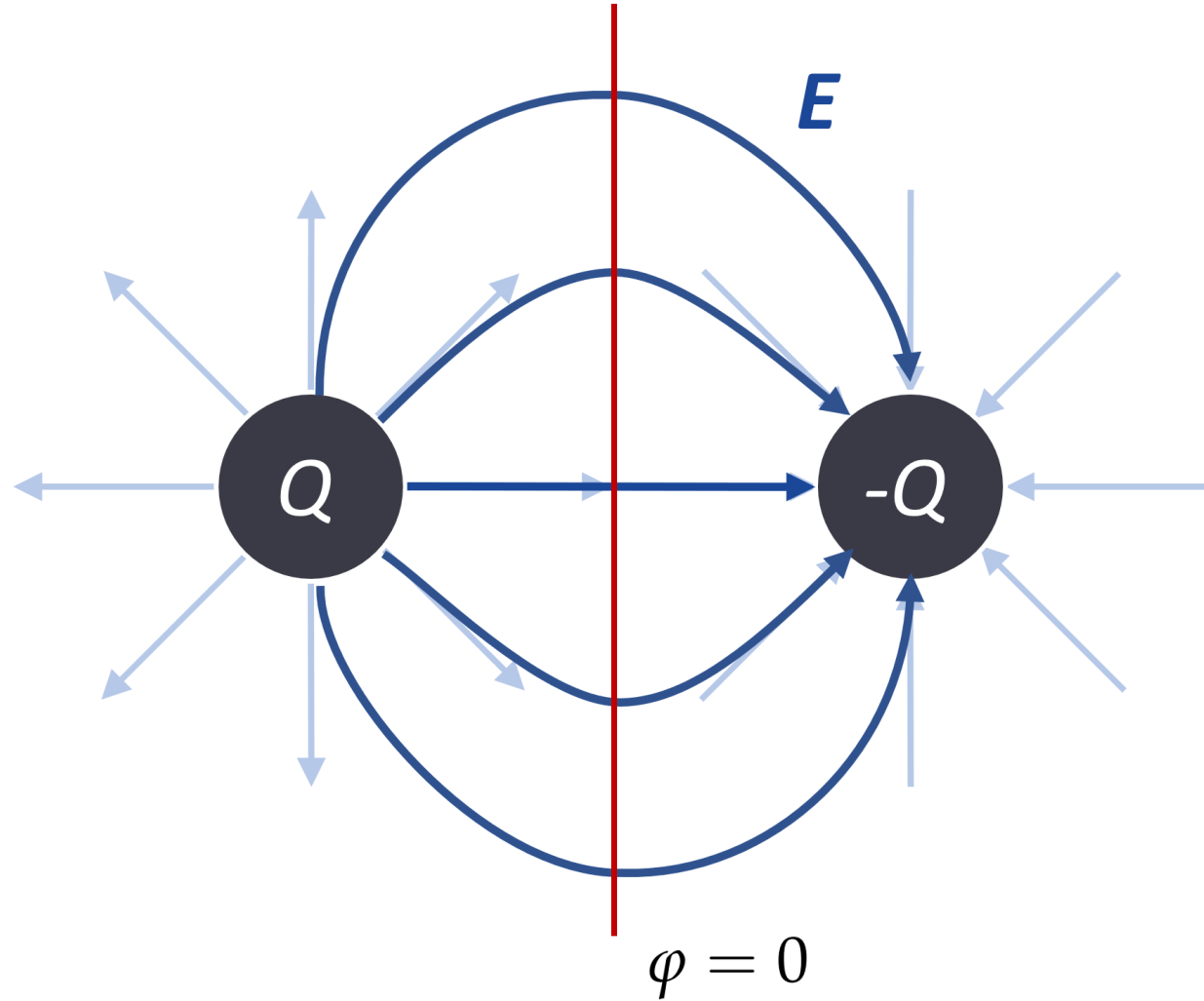
$$\mathbf{E}(\mathbf{r}) \sim \frac{1}{r^3}$$

$$\mathbf{p} = Q\mathbf{a}$$



[https://upload.wikimedia.org/wikipedia/commons/a/aa/VFPT\\_dipole\\_animation\\_electric.gif](https://upload.wikimedia.org/wikipedia/commons/a/aa/VFPT_dipole_animation_electric.gif)

# Electric dipole



# Electrodynamics

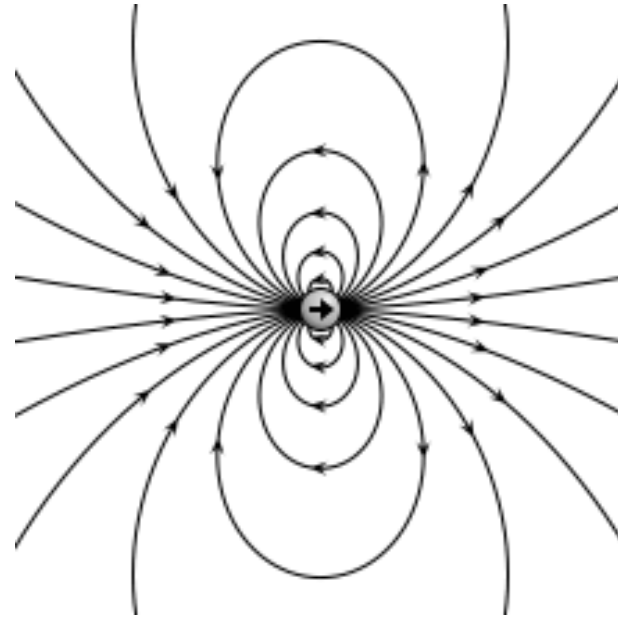
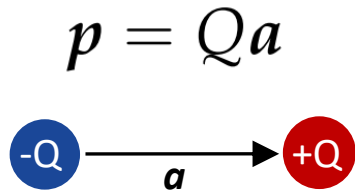
SECTION

Electrostatics

LECTURE

Force and torque acting  
on an electric dipole

# Electric dipole

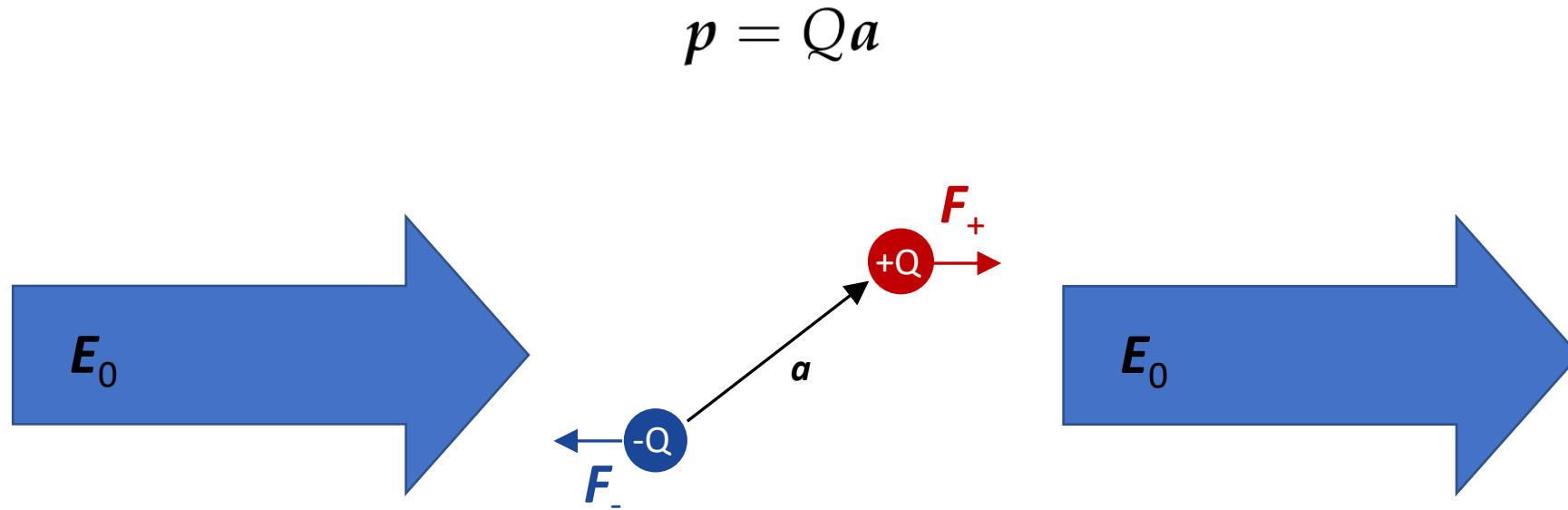


[https://upload.wikimedia.org/wikipedia/commons/a/aa/VFPT\\_dipole\\_animation\\_electric.gif](https://upload.wikimedia.org/wikipedia/commons/a/aa/VFPT_dipole_animation_electric.gif)

- Electric field

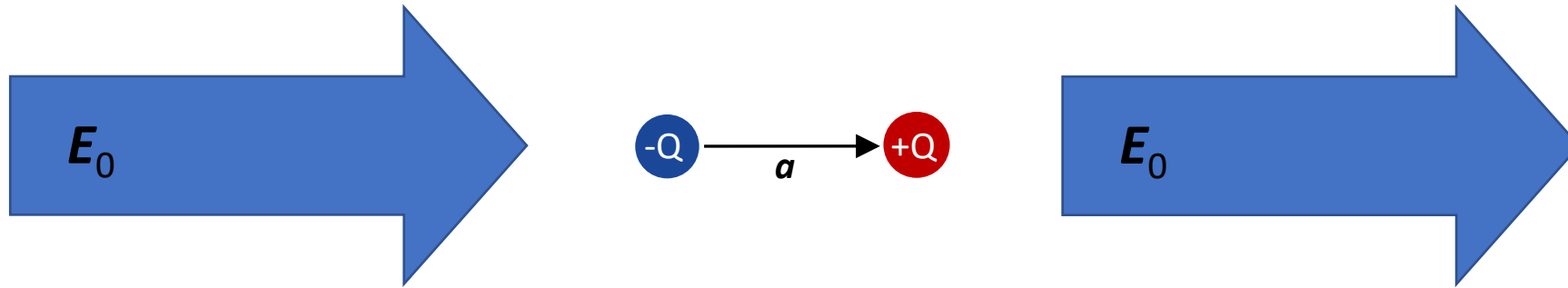
$$\mathbf{E}(\mathbf{r}) = -\frac{1}{4\pi\epsilon_0} \frac{r^2 \mathbf{p} - 3\mathbf{r}(\mathbf{p} \cdot \mathbf{r})}{r^5}$$

# Electric dipole in a homogeneous $E$ field



# Electric dipole in a homogeneous $E$ field

$$p = Qa$$

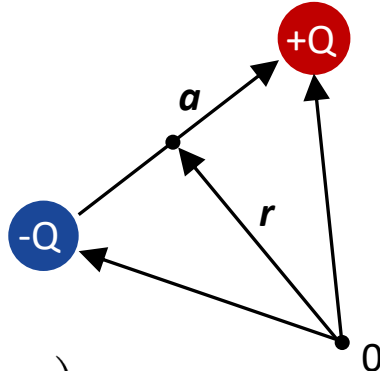


Reorientation  
No motion

# Electric dipole in a homogeneous $E$ field

**Force**

*No motion*



**Torque**

*Reorientation*

$$\mathbf{M} = \mathbf{r} \times \mathbf{F}$$

$$\mathbf{F} = Q\mathbf{E}(\mathbf{r} + 1/2\mathbf{a}) - Q\mathbf{E}(\mathbf{r} - 1/2\mathbf{a})$$

$$= Q[\mathbf{E}(\mathbf{r}) + (1/2\mathbf{a} \cdot \nabla)\mathbf{E}(\mathbf{r})] + \dots$$

$$- Q[\mathbf{E}(\mathbf{r}) - (1/2\mathbf{a} \cdot \nabla)\mathbf{E}(\mathbf{r})] + \dots$$

$$= Q(\mathbf{a} \cdot \nabla)\mathbf{E}(\mathbf{r}) + \dots$$

$$\mathbf{F} = (\mathbf{p} \cdot \nabla)\mathbf{E}(\mathbf{r})$$

$$\mathbf{F} = 0 \quad \text{for } \mathbf{E}(\mathbf{r}) = \mathbf{E}_0 = \text{const.}$$

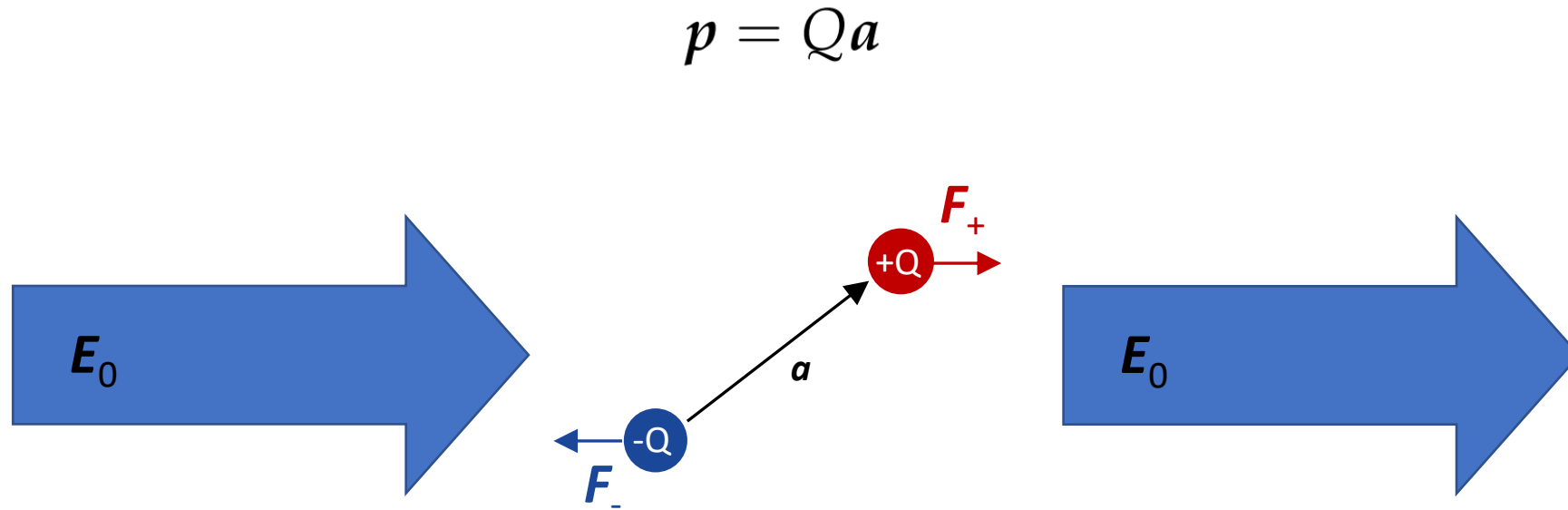
$$\mathbf{M} = 1/2\mathbf{a} \times Q\mathbf{E}(\mathbf{r} + 1/2\mathbf{a}) + 1/2\mathbf{a} \times Q\mathbf{E}(\mathbf{r} - 1/2\mathbf{a})$$

$$= 1/2\mathbf{a} \times Q\{[\mathbf{E}(\mathbf{r}) + (1/2\mathbf{a} \cdot \nabla)\mathbf{E}(\mathbf{r})] + \dots \\ + [\mathbf{E}(\mathbf{r}) - (1/2\mathbf{a} \cdot \nabla)\mathbf{E}(\mathbf{r})] + \dots\}$$

$$= Q\mathbf{a} \times \mathbf{E}(\mathbf{r}) + \dots$$

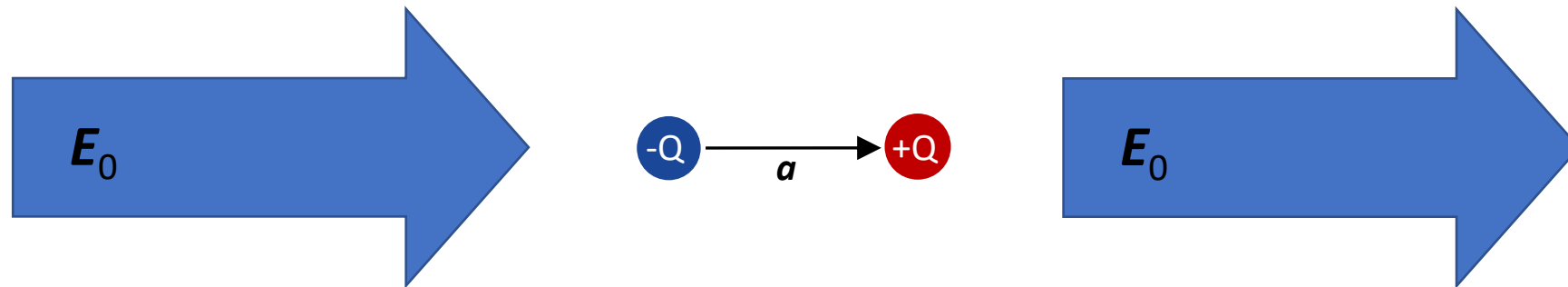
$$\mathbf{M} = \mathbf{p} \times \mathbf{E}(\mathbf{r})$$

# Electric dipole in a homogeneous $E$ field



# Electric dipole in a homogeneous $E$ field

$$\mathbf{p} = Q\mathbf{a}$$



Reorientation  $\mathbf{M} = \mathbf{p} \times \mathbf{E}(\mathbf{r})$

No motion  $\mathbf{F} = (\mathbf{p} \cdot \nabla)\mathbf{E}(\mathbf{r})$

# Electrodynamics

SECTION

Electrostatics

LECTURE

Boundary conditions  
& mirror charges

# Boundary conditions

$$\Delta\varphi(\mathbf{r}) = -\frac{1}{\epsilon_0}\rho(\mathbf{r}) \quad \text{Poisson equation}$$

- Special solution

$$\varphi_s(\mathbf{r}) = \int \frac{1}{4\pi\epsilon_0} \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'$$

- General solution

$$\varphi(\mathbf{r}) = \varphi_s(\mathbf{r}) + \varphi_h(\mathbf{r}) \quad \text{Solution of the homogeneous equation}$$

$$\Delta\varphi_h(\mathbf{r}) = 0$$

... has to fulfill boundary conditions

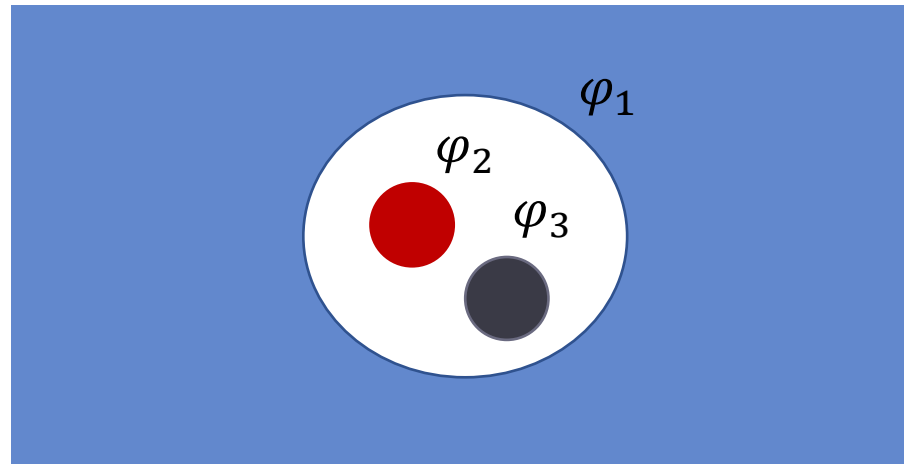
$$\text{so far } \varphi = 0 \quad \text{for } \mathbf{r} \rightarrow \infty \quad \varphi_h(\mathbf{r}) = 0$$

# Boundary conditions

$$\varphi(\mathbf{r}) = \varphi_s(\mathbf{r}) + \boxed{\varphi_h(\mathbf{r})} \quad \dots \text{has to fulfill boundary conditions}$$

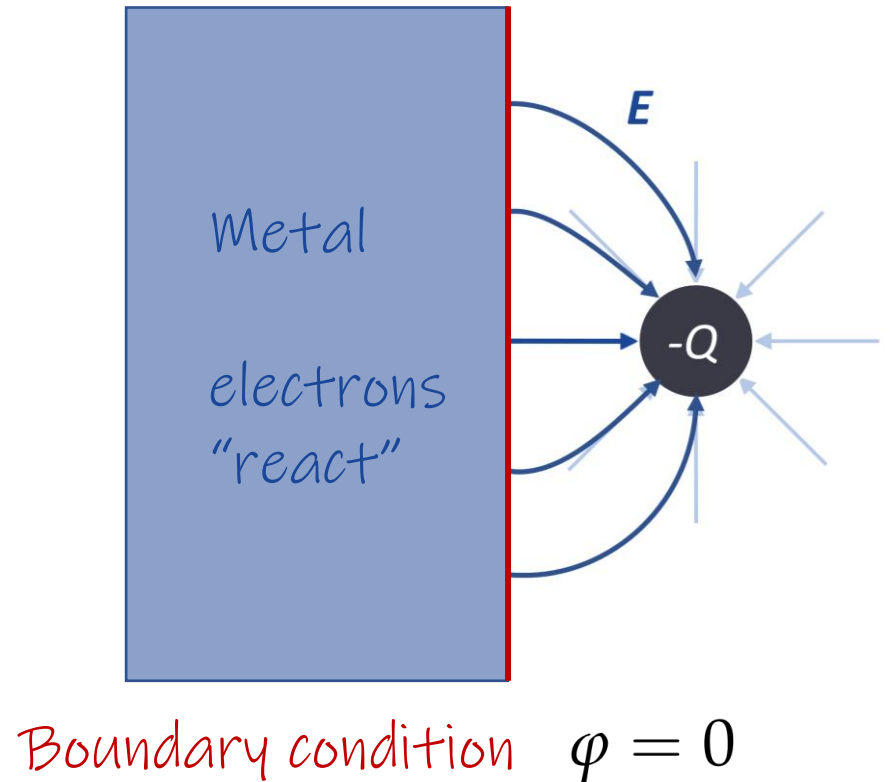
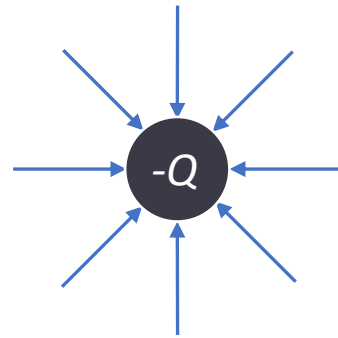
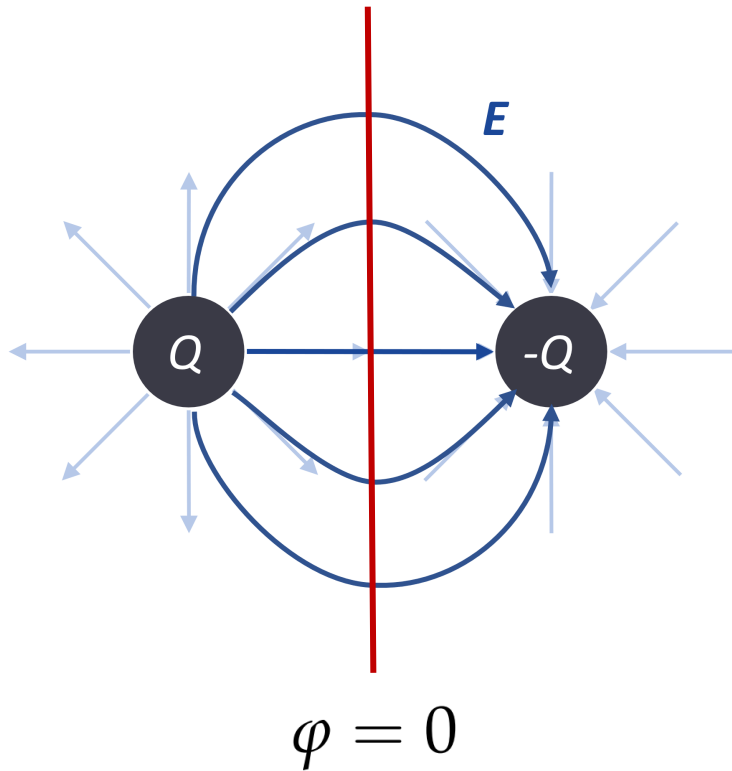
so far  $\varphi = 0$  for  $r \rightarrow \infty$   $\varphi_h(\mathbf{r}) = 0$

- Charged wire or metal surface: constant potential



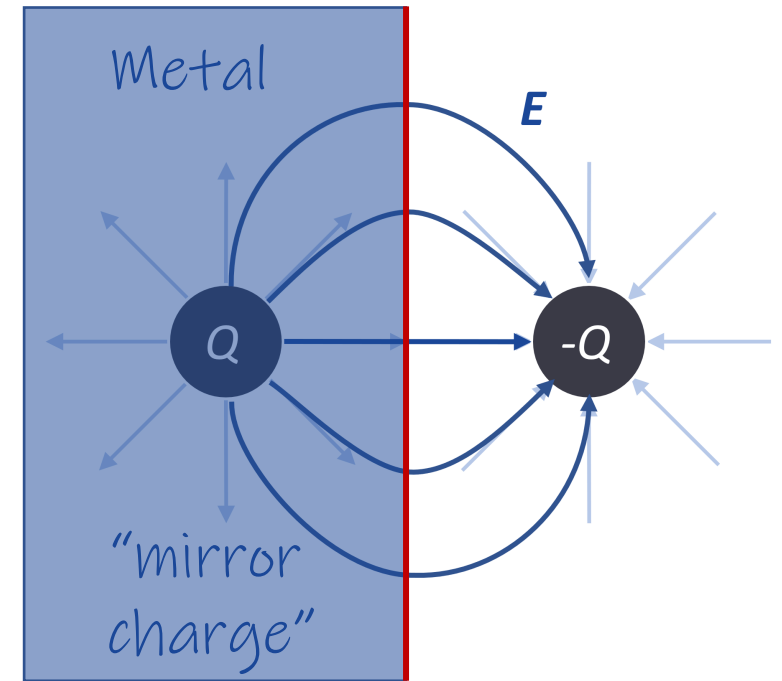
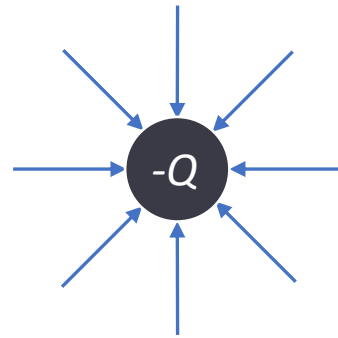
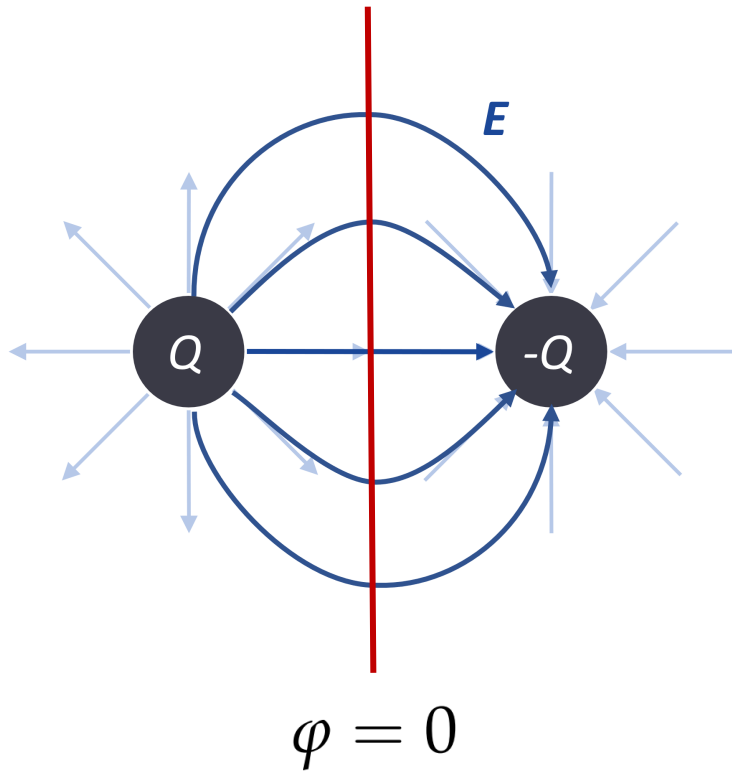
# Boundary conditions

$$\varphi(\mathbf{r}) = \varphi_s(\mathbf{r}) + \varphi_h(\mathbf{r}) \quad \dots \text{has to fulfill boundary conditions}$$



# Boundary conditions

$$\varphi(\mathbf{r}) = \varphi_s(\mathbf{r}) + \varphi_h(\mathbf{r}) \quad \dots \text{has to fulfill boundary conditions}$$



Boundary condition  $\varphi = 0$

# Electrodynamics

SECTION

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LECTURE

Energy in  
electrostatics

# Field energy

$$-\mathbf{j} \cdot \mathbf{E} = \frac{\partial}{\partial t} \left( \epsilon_0 \frac{\mathbf{E}^2}{2} + \frac{1}{2\mu_0} \mathbf{B}^2 \right) + \nabla \cdot \frac{\mathbf{E} \times \mathbf{B}}{\mu_0}$$

$$\nu = \frac{\partial}{\partial t} w + \nabla \cdot \mathbf{S}$$

Continuity equation for energy density

- Power density  
(energy generation density)

$$\nu = -\mathbf{j} \cdot \mathbf{E}$$

- Energy density

$$w = \epsilon_0 \frac{\mathbf{E}^2}{2} + \frac{1}{2\mu_0} \mathbf{B}^2$$

- Poynting vector  
(energy-current density)

$$\mathbf{S} = \frac{\mathbf{E} \times \mathbf{B}}{\mu_0}$$

# Electrostatic energy

- Energy density  $w = \epsilon_0 \frac{E^2}{2} + \frac{1}{2\mu_0} B^2$

- Energy  $W_e = \int w_e dV = \int \frac{\epsilon_0}{2} E^2 dV$

$$W_e = - \int \frac{\epsilon_0}{2} \mathbf{E} \cdot \nabla \phi dV$$

$$\mathbf{E}(\mathbf{r}) = -\nabla \phi(\mathbf{r})$$

$$W_e = - \int \frac{\epsilon_0}{2} \nabla \cdot (\mathbf{E} \phi) dV + \int \frac{\epsilon_0}{2} \phi \nabla \cdot \mathbf{E} dV$$

$$\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{1}{\epsilon_0} \rho(\mathbf{r})$$

$$W_e = - \oint \frac{\epsilon_0}{2} (\mathbf{E} \phi) \cdot d\mathbf{S} + \frac{1}{2} \int \rho \phi dV$$

$$W_e = \frac{1}{2} \int \rho(\mathbf{r}) \phi(\mathbf{r}) dV$$

If no surface charges

# Electrostatic energy

$$W_e = \int w_e dV = \int \frac{\epsilon_0}{2} E^2 dV$$

$$W_e = \frac{1}{2} \int \rho(\mathbf{r}) \varphi(\mathbf{r}) dV$$

If no surface charges

$$\varphi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} dV'$$

$$W_e = \frac{1}{8\pi\epsilon_0} \int \int \frac{\rho(\mathbf{r}')\rho(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|} dV dV' \quad \text{general}$$

$$W_e = \frac{1}{8\pi\epsilon_0} \sum_{i \neq j} \frac{Q_i Q_j}{|\mathbf{r}_i - \mathbf{r}_j|} \quad \text{point charges}$$

$$W_e = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{|\mathbf{r}_1 - \mathbf{r}_2|} \quad \text{2 charges}$$

# Electrodynamics

SECTION

Electrostatics

LECTURE

Summary

# Electrostatics

$$\nabla \times \mathbf{E}(\mathbf{r}) = 0$$

$$\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{1}{\epsilon_0} \rho(\mathbf{r})$$

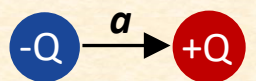
$$\mathbf{E}(\mathbf{r}) = -\nabla \phi(\mathbf{r})$$

$$\Delta \phi(\mathbf{r}) = -\frac{1}{\epsilon_0} \rho(\mathbf{r})$$

$$\phi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} dV'$$

$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \rho(\mathbf{r}') \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3} dV'$$

$\mathbf{p} = Q\mathbf{a}$



$$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{3(\mathbf{p} \cdot \mathbf{r})\mathbf{r} - p r^2}{r^5}$$

$$\mathbf{F} = (\mathbf{p} \cdot \nabla)\mathbf{E}(\mathbf{r}) \quad \mathbf{M} = \mathbf{p} \times \mathbf{E}(\mathbf{r})$$

$$W_e = \int w_e dV = \int \frac{\epsilon_0}{2} \mathbf{E}^2 dV$$

$$W_e = \frac{1}{2} \int \rho(\mathbf{r}) \phi(\mathbf{r}) dV$$

